

# Heteroepitaxial Growth of Nanostructured Cerium Dioxide Thin Films by MOCVD on a (001) TiO<sub>2</sub> Substrate

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The deposition of epitaxial CeO<sub>2</sub> nanostructured thin films on rutile (TiO<sub>2</sub>) substrates by metal–organic chemical vapor deposition (MOCVD) has been carried out in a wide temperature range. Films have been grown on (001)TiO<sub>2</sub> from the Ce(III) 1,1,1,5,5,5-hexafluoro-2,4-pantanediolato diglyme adduct (Ce(hfa)<sub>3</sub>·diglyme). The X-ray diffraction patterns of all samples grown in the 450–750 °C deposition temperature range point to the formation of {100}-oriented CeO<sub>2</sub> films, whereas at higher deposition temperatures (850–1050 °C) all the deposited CeO<sub>2</sub> films show {111} texture. fwhm rocking curves values clearly indicate that the texturing of CeO<sub>2</sub> crystallites is greatly improved upon increasing the deposition temperature in both cases. Typical pole figure patterns demonstrate a cube-on-cube epitaxial growth in the case of {100}-oriented CeO<sub>2</sub> films, whereas {111}-oriented CeO<sub>2</sub> films show different in-plane alignments of the crystallites. A possible explanation for textural changes is proposed.

## Introduction

Major advances have been made in recent years in high-temperature superconductor (HTS) research, resulting in an increasing use of HTS materials in commercial and precommercial power applications<sup>1–3</sup> and electronic devices.<sup>4,5</sup> In particular, HTS films are good candidates for devices adopted in cellular communication areas.<sup>6–9</sup> In this context, rutile (TiO<sub>2</sub>) has attracted considerable interest in recent years as a versatile and suitable substrate for fabrication of HTS passive filters. In fact, this substrate possesses a higher

dielectric constant with respect to LaAlO<sub>3</sub>, YSZ, and Al<sub>2</sub>O<sub>3</sub> and relatively low loss tangent.<sup>10</sup> The most important point in these applications is the preparation of highly (001)-oriented HTS thin films to achieve superconducting devices with low microwave losses.<sup>11,12</sup>

In the case of HTS applications for electronic and magnetic devices, texturing plays a key role since well-matched and -textured multilayer structures often drive properties and performances as well as increase the device lifetime. Moreover, the growth of epitaxial films currently provides the opportunity to study film and multilayer properties that depend on the orientation of crystalline growth.

Finally, the fabrication of HTS multilayer devices requires lattice-matched, chemically compatible, and thermally stable materials as buffer, dielectric, and/or insulator layers.

Among the numerous oxides often used as buffer layers,<sup>13</sup> epitaxial cerium oxide (CeO<sub>2</sub>) thin films are currently of interest for different applications in key areas of thin film technology, ranging from semiconductors<sup>14–16</sup> to superconductors<sup>17,18</sup> and solid oxide fuel cells.<sup>19</sup> In fact, these films are used as protective buffer

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(1) Wu, X. D.; Foltyn, S. R.; Arendt, P. N.; Blumenthal, W. R.; Campbell, I. H.; Cotton, J. D.; Coulter, J. Y.; Hults, W. L.; Maley, M. P.; Safar, H. F.; Smith, J. L. *Appl. Phys. Lett.* **1995**, *67*, 2397–2399.

(2) Goyal, A.; Norton, D. P.; Budai, J. D.; Paranthaman, M.; Specht, E. D.; Koerger, D. M.; Christen, D. K.; He, Q.; Saffian, B.; List, F. A.; Lee, D. F.; Martin, P. M.; Klabund, C. E.; Hartfield E.; Sikka, V. K. *Appl. Phys. Lett.* **1996**, *69*, 1795–1797.

(3) Matsumoto, K.; Kim, S. B.; Win, J. G.; Hirabayashi, I.; Watanabe, T.; Unp, N.; Ikeda, M. *IEEE Trans. Appl. Supercond.* **1999**, *9*, 1539–1542.

(4) Terashima, Y.; Yamazaki, M.; Kudo, Y.; Yoshino, H.; Ando, K.; Oshima, S. *Jpn. J. Appl. Phys.* **1994**, *33*, L1592–L1594.

(5) Chaloupka, H. J.; Hein, M. A.; Muller, G. *SPIE Proc.* **1994**, *36*, 2156.

(6) Lancaster, M. J. *Passive Microwave Device Applications of High-Temperature Superconductors*; Cambridge University Press: Cambridge, 1997; Chapters 4 and 5.

(7) Hein, M. *High-Temperature-Superconductor Thin Films at Microwave Frequencies*; Springer Tracts in Modern Physics; Springer: Berlin, 1999; Vol. 155, Chapters 3 and 5.

(8) Klaudam, M.; Kasser, T.; Mayer, B.; Neumann, C.; Schnell, F.; Aminov, B.; Baumfalk, A.; Chaloupka, H.; Kolesov, S.; Piel, H.; Klein, N.; Schornstein, S.; Bareiss, M. *IEEE Trans. Microwave Theory Tech.* **2000**, *48*, 1227–1230.

(9) Li, H.; He, Y. S.; He, A. S.; Li, S.; Li, C.; Yon, L.; Liang, J.; Zhu, W.; Zhou, Y.; Hong, J.; Lancaster, M. J. *Supercond. Sci. Technol.* **2002**, *15*, 276–279.

(10) Templeton, A.; Wang, X. R.; Penn, S. J.; Webb, S. J.; Cohen, L. F.; Alford, N. M. *J. Am. Ceram. Soc.* **2000**, *83*, 95–98.

(11) Nerumi, E.; Song, L. W. *Appl. Phys. Lett.* **1991**, *58*, 1202–1204.

(12) Willensen, B. A.; Kilhstrom, K. E.; Dahlen, T.; Scalapino, D. J.; Gowe, B.; Bonn, D. A.; Hardy, W. N. *Phys. Rev. B* **1998**, *58*, 6650–6654.

(13) Wordenweber, R. *Supercond. Sci. Technol.* **1999**, *12*, R86–R102.

(14) Cossarutto, L.; Chaoui, N.; Million, E.; Muller, J. F.; Lambert, J.; Alnot, M. *Appl. Surf. Sci.* **1998**, *126*, 352–355.

(15) Jones, J. T.; Bridger, P. M.; Marsh, O. J.; McGill, T. C. *Appl. Phys. Lett.* **1999**, *75*, 1326–1328.

(16) Baue, G.; Song, Y.; Jung, D.; Roh, Y. *Appl. Phys. Lett.* **2000**, *77*, 729–731.

layers for HTS growth, owing to the small lattice mismatch with 45° rotated crystallites of HTS films, to similar thermal expansion coefficients,<sup>20</sup> and to their chemical stability.

$\text{CeO}_2$  thin films have been deposited by several techniques including thermal and electron beam evaporation,<sup>21,22</sup> laser ablation,<sup>23</sup> pulsed laser deposition,<sup>24</sup> ion-assisted deposition,<sup>25</sup> and reactive dc magnetron sputtering.<sup>26</sup>

Metal-organic chemical vapor deposition (MOCVD), however, represents a more attractive technique for microelectronic applications since it offers great potentialities for large area growth, good film uniformity, and excellent conformal step coverage at device dimensions less than 2  $\mu\text{m}$ . The majority of MOCVD studies, focused on fabrication of  $\text{CeO}_2$  thin films with superior crystalline and morphological quality, have used single-crystal substrates, which show a little lattice mismatch relative to  $\text{CeO}_2$  films, such as yttria-stabilized zirconia (YSZ) (5%),<sup>27,28</sup>  $\text{Al}_2\text{O}_3$  r-plane (5.7% and 12%),<sup>20,29</sup> and Si (0.4%).<sup>30</sup> Generally, for substrates with a reasonable lattice match, operational conditions required for film depositions are relatively substrate-independent.  $\text{CeO}_2$  and  $\text{TiO}_2$  belong to different space groups. The former has a cubic fluorite-type structure with a lattice parameter  $a = 0.5411$  nm, whereas the latter belongs to a tetragonal structure with lattice spacings  $a = 0.4593$  nm and  $c = 0.2959$  nm. In this case, (001)  $\text{TiO}_2$  substrate and (100)  $\text{CeO}_2$  thin films are largely mismatched (about 17%).<sup>31</sup> This large difference points to a more critical epitaxial growth than that observed on previously mentioned substrates.

The present study is addressed to find the best deposition conditions for fabrication of high-quality  $\text{CeO}_2$  buffer layers on  $\text{TiO}_2$ . To our knowledge, there have been no previous reports on the growth of  $\text{CeO}_2$  on  $\text{TiO}_2$  substrates.

$\text{CeO}_2$  thin films have been grown by MOCVD using a volatile metal-organic second-generation precursor,  $\text{Ce}(\text{hfa})_3$ ·diglyme [ $\text{Hhfa} = 1,1,1,5,5,5$ -hexafluoro-2,4-pentanedione, diglyme =  $(\text{CH}_3\text{O}(\text{CH}_2\text{CH}_2\text{O})_2\text{CH}_3)$ ]. The

epitaxy of (100)-oriented  $\text{CeO}_2$  films has been driven by varying deposition parameters. It has been found that nanocrystalline  $\text{CeO}_2$  thin films grow with different crystallographic orientations on rutile substrates upon varying the deposition temperature.

## Experimental Section

A horizontal hot-wall MOCVD reactor was used in all experiments. The anhydrous, air, and thermally stable metallorganic source,  $\text{Ce}(\text{hfa})_3$ ·diglyme, was prepared as described elsewhere.<sup>32</sup> It was vaporized from the melt at 120 °C. Ar (100 sccm) was used as a carrier gas and 100 sccm water-saturated oxygen flow was fed close to the substrate as the reacting gas. The water bubbler was kept at 25 °C. Fluorine contamination in the films was precluded by adding water vapor into the oxygen stream.<sup>33</sup> Depositions were carried out for 60 min in the reactor chamber maintained at 4 Torr. Deposition temperatures were varied in the 450–1050 °C range. The deposition rate changed from 1.1  $\mu\text{m}/\text{h}$  at 450 °C deposition temperature to 0.7  $\mu\text{m}/\text{h}$  at 1050 °C. The thicknesses were evaluated by cross-section SEM observations. (001)-oriented  $\text{TiO}_2$  10  $\times$  10 mm<sup>2</sup> substrates were purchased from Crystal GmbH.

X-ray diffraction measurements (XRD) were recorded on a Bruker-AXS D5005  $\theta$ – $\theta$  X-ray diffractometer, using Cu K $\alpha$  radiation operating at 40 kV and 30 mA. Rocking curves of the  $\text{CeO}_2$  (200) and (111) reflections were measured and used as texturing indicators. A four circle Bruker-AXS D5005 X-ray diffractometer was used to study the  $\text{CeO}_2$  films in-plane grain alignment.

Film surface morphologies were examined using a LEO Iridium 1450 scanning electron microscope (SEM). Film atomic composition was determined by wavelength dispersive X-ray analysis (WDX) using an Oxford Instruments WDX-3PC analyzer. The WDX analyses were carried out at 15-kV accelerating voltage, 10-nA current, and 100-s acquisition time. TAP (thallium acid phthalate) and LSM60 (synthetic multi-layered structure) crystals were used as an X-ray selector for F and C, respectively.

## Results and Discussion

**Structural Characterization.** Earlier results<sup>32,34</sup> have already shown that the presently adopted  $\text{Ce}(\text{hfa})_3$ ·diglyme precursor represents a reliable source for deposition of cerium-containing films due to its low melting point (75 °C), which allows evaporation from the liquid phase. This represents an important issue for MOCVD applications, where liquid source precursors appear the most desirable due to the greater stability of the vapor pressure and to the absence of any effect of crystallite size on the evaporation rate and, hence, on the film growth rate. Even more important,  $\text{Ce}(\text{hfa})_3$ ·diglyme adduct produces good quality  $\text{CeO}_2$  thin films at a deposition temperature (450 °C) significantly lower than that usually required (>550 °C).<sup>35</sup> This also is an important issue since lower deposition temperatures are highly desirable for multilayer device fabrication.

Initial experiments for growing  $\text{CeO}_2$  films over  $\text{TiO}_2$  (001) substrates have been attempted at 450 °C. The substrate size is 10  $\times$  10 mm<sup>2</sup> and the films are uniform

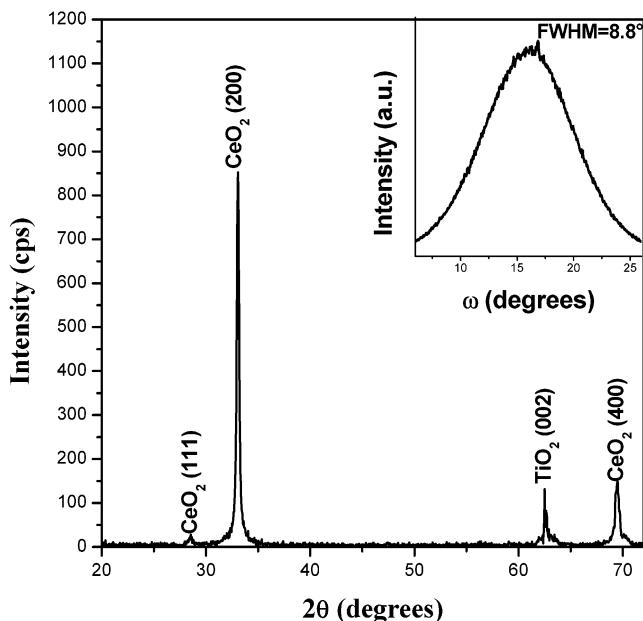
- (17) Inoue, T.; Yamamoto, Y.; Koyama, S.; Suzuki, S.; Ueda, Y. *Appl. Phys. Lett.* **1990**, *56*, 1332–1333.
- (18) Tian, Y. J.; Linzen, S.; Schmidl, F.; Matthes, A.; Scheelidewind, H.; Siedel, P. *Thin Solid Films* **1999**, *338*, 224–230.
- (19) Inoue, T.; Setoguchi, T.; Educhi, K.; Arai, H. *Solid State Ion.* **1989**, *35*, 285–291.
- (20) Lu, Z.; Hiskes, R.; Di Carolis, S. A.; Nel, A.; Rate, R. K.; Feigelson, R. S. *J. Cryst. Growth* **1995**, *156*, 227–234.
- (21) Lee, S.; Oh, D.; Goo, D.; Youm, D. *Thin Solid Films* **1995**, *258*, 299–304.
- (22) Inoue, T.; Ohsuna, T. *Appl. Phys. Lett.* **1991**, *59*, 3604–3606.
- (23) Yoshimoto, M.; Nagata, H.; Tsukahara, T.; Koinuma, H. *Jpn. J. Appl. Phys.* **1990**, *29*, L1199–L1202.
- (24) Nagata, H.; Tsukahara, T.; Gonda, A.; Yoshimoto, M.; Koinuma, H. *Jpn. J. Appl. Phys.* **1991**, *30*, L1136–L1138.
- (25) Al-Robaee, M. S.; Naraimha, K.; Mohan, S. *J. Appl. Phys.* **1992**, *71*, 2380–2386.
- (26) Granqvist, C. G.; Azens, A.; Hjelm, A.; Kullman, L.; Niklasson, G. A.; Ronnow, D.; Stromme Mattsson, M.; Veszelei, M.; Vaisar, G. *Sol. Energy* **1998**, *63*, 199–216.
- (27) Becht, M.; Morishita, T. *J. Alloys Compd.* **1997**, *251*, 310–313.
- (28) Belot, J. A.; Wang, A. C.; McNeely, R. J.; Liable-Sanfs, L.; Rheingold, A. L.; Marks, T. J. *Chem. Vap. Dep.* **1999**, *5*, 65–69.
- (29) Frohlich, K.; Souc, J.; Machajdik, D.; Jergel, M.; Snaauwaert, J.; Hellemans, L. *Chem. Vap. Deposition* **1998**, *6*, 216–220.
- (30) Fukuda, H.; Miura, M.; Sakuma, S.; Nomura, S. *Jpn. J. Appl. Phys., Part 1* **1998**, *37*, 4158–4159.
- (31) The following notation has been used:  $[hkl]$  denotes a single direction and  $\langle hkl \rangle$  a general set of directions; a plane is labeled as  $(hkl)$  and a general set of planes as  $\{hkl\}$ .

(32) Malandrino, G.; Lo Nigro, R.; Benelli, C.; Castelli, F.; Fragalà, I. L. *Chem. Vap. Deposition* **2000**, *6*, 233–238.

(33) McAleese, J.; Plakatouras, J. C.; Steele, B. C. H. *Thin Solid Films* **1996**, *280*, 152–159.

(34) Lo Nigro, R.; Malandrino, G.; Fragalà, I. L. *Chem. Mater.* **2001**, *13*, 4402–4404.

(35) Wang, A. C.; Belot, J. A.; Marks, T. J.; Markworth, P. R.; Chang, R. P. H.; Chudzik, M. P.; Kannewurf, C. R. *Physica C* **1999**, *320*, 154–160.

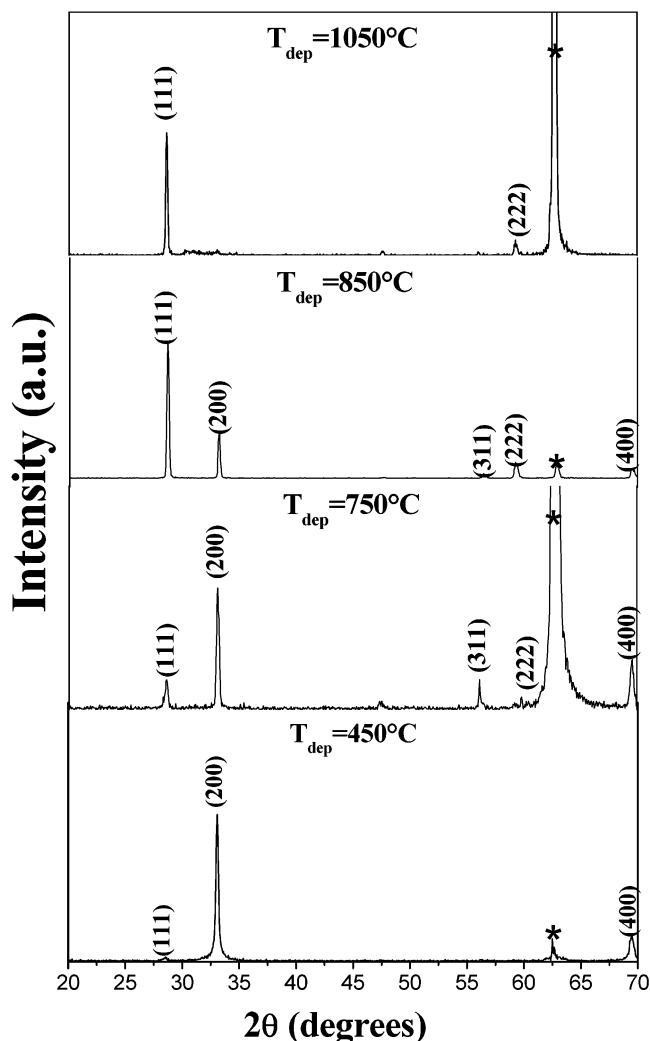


**Figure 1.**  $\theta$ - $2\theta$  XRD scan of a (100)  $\text{CeO}_2$  film deposited at 450  $^{\circ}\text{C}$  on  $\text{TiO}_2$  (001). The insert shows the rocking curve of the (200) reflection.

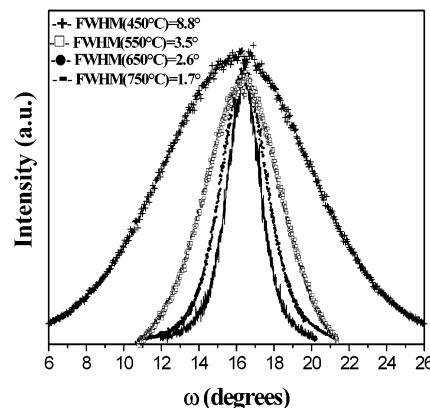
all over the surface. The XRD pattern is reported in Figure 1. The presence of peaks at  $2\theta = 33.08^{\circ}$  and  $69.4^{\circ}$ , related to the (200) and (400) reflections, respectively, indicates that  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  thin films have been obtained on  $\text{TiO}_2$  (001) at 450  $^{\circ}\text{C}$ . The rocking curve (see insert in Figure 1) of the (200)  $\text{CeO}_2$  reflection has a full-width at half-maximum (fwhm) of about 9°, significantly greater than that found for  $\text{CeO}_2$  films deposited on single-crystal substrates at the same deposition temperature.<sup>27,32,35</sup> This effect certainly accompanies the great lattice mismatch with (001)  $\text{TiO}_2$  substrates.

To improve epitaxy, several depositions have been attempted in the 450–1050  $^{\circ}\text{C}$  substrate temperature range and the consequent texture evolution has been monitored. In the 450–750  $^{\circ}\text{C}$  temperature range,  $\text{CeO}_2$  films are mostly  $\langle 100 \rangle$ -oriented with an evident gradual change to  $\langle 111 \rangle$  texturing upon increasing the temperature (Figure 2). In the same range, the fwhm values of the (200) reflection rocking curves (Figure 3) decrease, thus clearly showing that the alignment of  $\text{CeO}_2$  films improves significantly. Samples deposited at 750  $^{\circ}\text{C}$  are  $\langle 100 \rangle$ -oriented and show a considerably smaller fwhm value (1.7°). These films contain small quantities of (111)-oriented grains. This represents a recurring problem in the preparation of  $\text{CeO}_2$  films by various techniques<sup>36</sup> since the competing growth of  $\langle 111 \rangle$ -oriented domains depends on the deposition temperature. Nevertheless, a small quantity of  $\langle 111 \rangle$  grains in  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  films is not a critical issue for their use as buffer layers.

Further investigations above 750  $^{\circ}\text{C}$  have shown an intriguing competition between  $\langle 100 \rangle$  and  $\langle 111 \rangle$  domains. Films deposited at 850  $^{\circ}\text{C}$  are mostly  $\langle 111 \rangle$ -textured, even though there is evidence of some  $\langle 100 \rangle$  crystals. The (111) domains become predominant at 1050  $^{\circ}\text{C}$ . Representative  $\theta$ - $2\theta$  scans are reported in Figure 2.



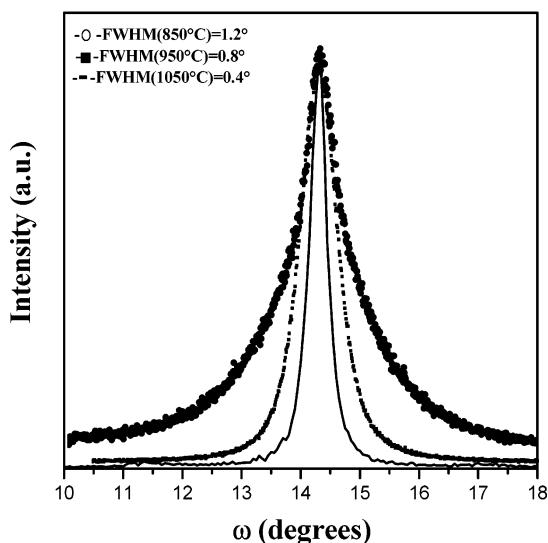
**Figure 2.**  $\theta$ - $2\theta$  XRD patterns of  $\text{CeO}_2$  films deposited at various temperatures (450–1050  $^{\circ}\text{C}$ ) (\* = substrate (002) reflection).



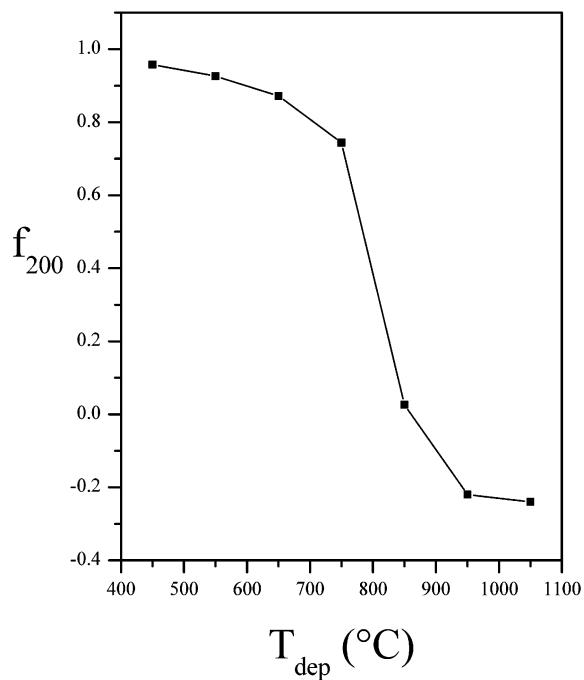
**Figure 3.** Rocking curves of the (200) reflection of  $\text{CeO}_2$  films at various deposition temperatures (450–750  $^{\circ}\text{C}$ ).

Rocking curves of the (111) reflection at selected temperatures are shown in Figure 4. There is clear evidence that narrower widths are associated with films deposited at higher temperatures and the smallest fwhm value (0.4°) has been measured at 1050  $^{\circ}\text{C}$ . fwhms measured on both  $\langle 100 \rangle$  and  $\langle 111 \rangle$  films generally point to a better texturing in the case of higher deposition temperature. In particular, the  $\langle 111 \rangle$   $\text{CeO}_2$  films on  $\text{TiO}_2$  substrates, prepared at deposition temperature as high

(36) Frohlich, K.; Souc, J.; Machajdik, D.; Kobzey, A. P.; Weiss, F.; Senateur, J. P.; Dahmen, K. H. *J. Phys. IV* **1995**, 5, C5-533–C5-539.



**Figure 4.** Rocking curves of the (111) reflection of  $\text{CeO}_2$  films at various deposition temperatures (850–1050 °C).



**Figure 5.** Variation of the degree of (200) preferred orientation ( $f_{200}$ ) versus deposition temperatures.

as 1050 °C, show an out-of-plane alignment better than  $\langle 100 \rangle$   $\text{CeO}_2$  films.

The amount of  $\langle 100 \rangle$  relative to  $\langle 111 \rangle$  domains can be semiquantitatively evaluated from the plot of the  $f_{200}$  parameter (degree of  $\langle 100 \rangle$  preferred orientation) vs the temperature (Figure 5),

$$f_{200} = (P_{200} - P_{200}^0) / (1 - P_{200}^0)$$

$$P_{200} = I_{200} / (I_{111} + I_{200})$$

$$P_{200}^0 = I_{200}^0 / (I_{111}^0 + I_{200}^0)$$

where  $I_{hkl}$  and  $I_{hkl}^0$  are the intensities of  $(hkl)$  reflections in the present  $\text{CeO}_2$  films and in a polycrystalline standard powder sample, respectively. Of course, only the (111) and (200) peaks have been considered since

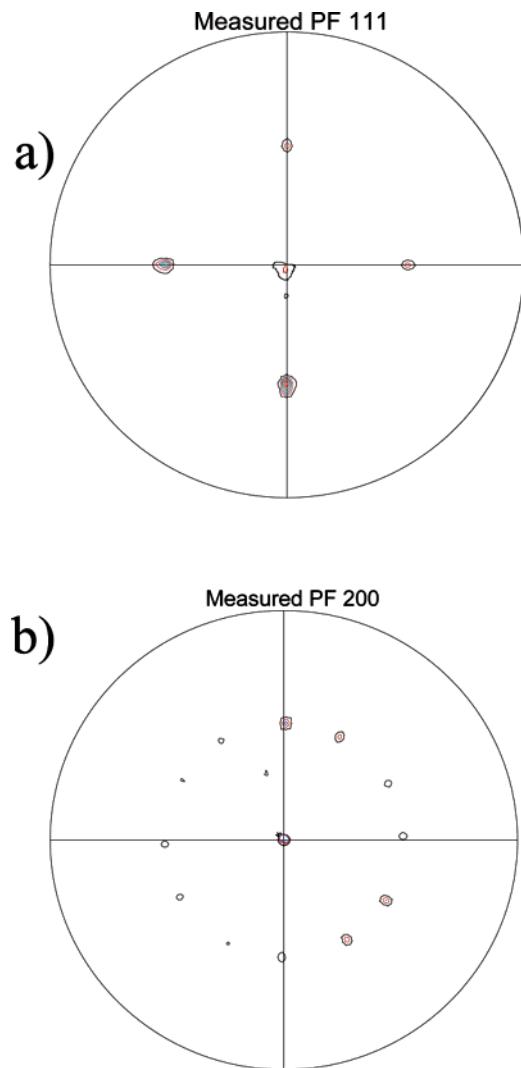
the other peaks have really negligible intensities. Negative values (as in the 900–1050 °C range) of the  $f_{200}$  parameter indicate that films are  $\langle 111 \rangle$  preferentially oriented. At lower deposition temperature (450 °C)  $\text{CeO}_2$  films are  $\langle 100 \rangle$ -oriented (i.e.,  $f_{200} \approx 100\%$ ). The rocking curve fwhm value points, however, to a poor crystallite alignment. Films deposited at 750 °C are 80%  $\langle 100 \rangle$ -textured with 1.7° fwhm. In the films deposited at 850 °C, the intensity of the (200) reflection is less than 30%, and the grain dispersion is so high that the rocking curve is no more detectable. A further increase of the deposition temperature (950 °C) results in a marked improvement of both  $\langle 111 \rangle$  texture (95%) and of the out-of-plane crystallite alignment (fwhm = 0.8°). At 1050 °C  $\text{CeO}_2$  films are 100%  $\langle 111 \rangle$ -oriented and present the best crystallite alignment (fwhm = 0.4°).

In regard to the film thicknesses, note that the deposition rate decreases on increasing temperature. This is a well-known issue since precursor depletion occurs at higher temperature due to its decomposition on the reactor walls. Nevertheless, there is no doubt that different growth rates might also depend on the different orientations.

**In-Plane Texture Analysis.** In this context, it becomes of interest to understand the epitaxy of  $\text{CeO}_2$  films and, in particular, whether the heterostructure  $\text{CeO}_2/\text{TiO}_2$  is in-plane-oriented. The pole figures, plotted in the 0°(center)  $< \psi < 90^\circ$  (rim) range of the distance of the pole angle, are reported in parts a and b of Figure 6. Figure 6a shows the pattern of  $\langle 100 \rangle$   $\text{CeO}_2$  thin films deposited at 750 °C. The  $2\theta$  angle has been fixed at 28.54° (corresponding to  $\text{CeO}_2$  (111) reflection) to record the pole figure. Four diffraction peaks, 90° spaced in  $\varphi$ , have been observed at  $\psi = 54.5^\circ$  with a 4-fold symmetry, showing, as expected for a cubic structure, that the  $\text{CeO}_2$  {111} planes are 54.75° tilted with respect to the  $\text{CeO}_2$  {100} planes, which lie parallel to the substrate surface. The four peaks, 90° spaced, clearly point to a good in-plane orientation. The central peak corresponds to the  $\langle 111 \rangle$ -oriented grains as observed in the X-ray pattern (Figure 2).

The epitaxial relationship between the  $\langle 111 \rangle$ -oriented  $\text{CeO}_2$  films and  $\langle 001 \rangle$   $\text{TiO}_2$  substrate has been studied as well. In the case of  $\langle 111 \rangle$ -oriented  $\text{CeO}_2$  thin films, deposited at 1050 °C, the related pole figure has been recorded using a  $2\theta$  angle of 33.08°, corresponding to the (200) reflection. Twelve equally spaced (every 30°) peaks at  $\psi = 54.6^\circ$  and one peak at  $\psi = 0^\circ$  have been observed, instead of the three peaks expected for the  $\langle 100 \rangle$  directions. This indicates that the  $\langle 111 \rangle$   $\text{CeO}_2$  films consist of, at least, two sets (vide infra) of crystallites, each rotated by 90° in the substrate plane. In addition, crystallites are probably twinned about the [111] direction, as expected for a face-centered cubic crystal.<sup>37</sup> On the basis of the cubic symmetry of this lattice, crystal growth could even involve four sets of twinned crystallites with the consequence that, for more than two crystallite in-plane orientations, various X-ray reflections could coincide. An obvious consequence, therefore, is that the exact number of distinct crystallite orientations cannot be determined. In Figure 7 a schematic

(37) Kelly, A.; Groves, G. W. *Crystallography and Crystal Defects*; Addison-Wesley: Reading, MA, 1970.

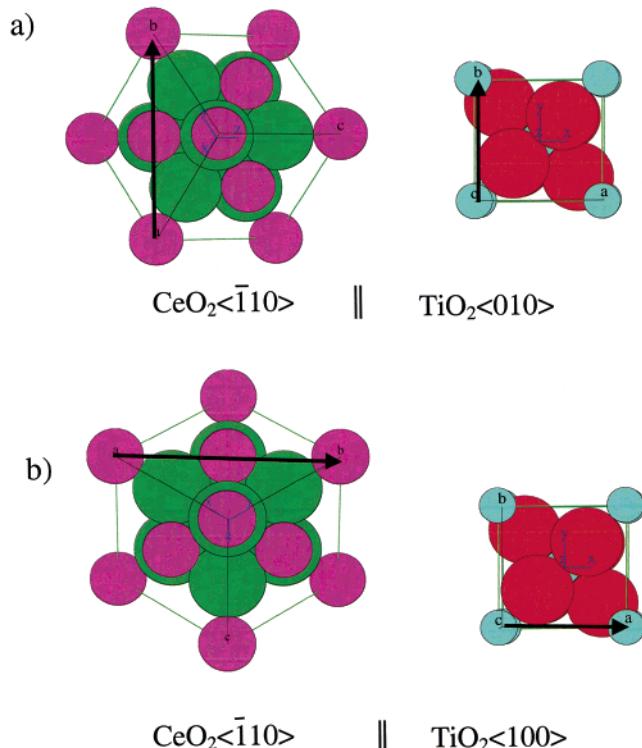


**Figure 6.** Pole figures of (a)  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  films deposited at  $750\text{ }^\circ\text{C}$ , using the (111) reflection as the pole; (b)  $\langle 111 \rangle$ -oriented  $\text{CeO}_2$  films deposited at  $1050\text{ }^\circ\text{C}$ , using the (200) reflection as the pole.

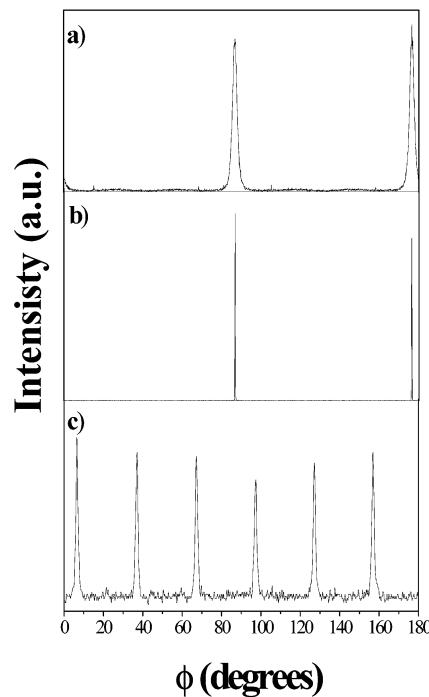
representation of the  $\text{CeO}_2$  (111) and  $\text{TiO}_2$  (001) planes is reported. If we consider the  $\text{CeO}_2$   $[-110]$  direction, which lies in the (111) plane, two epitaxial in-plane variants are possible. One alignment arises when the  $\text{CeO}_2$   $[-110]$  direction is parallel to the  $\text{TiO}_2$   $[010]$  (Figure 7a) and the second derives by a  $90^\circ$   $\text{CeO}_2$  cell rotation around the [111], with a consequent alignment of the  $\text{CeO}_2$   $[-110]$  with the  $\text{TiO}_2$   $[100]$  (Figure 7b).

Identical behavior has already been reported for other heteroepitaxial pairs, such as  $(111)\text{BaTiO}_3 \parallel (001)\text{YSZ}^{38}$  or  $(111)\text{LaAlO}_3 \parallel (001)\text{YSZ}^{39}$ .

Additional information about the in-plane relationship of  $\text{CeO}_2$  films relative to the underlying  $\text{TiO}_2$  substrate can be obtained from  $\Phi$ -scans of the  $\text{CeO}_2/\text{TiO}_2$  film/substrate system (Figure 8).  $\Phi$ -scans of  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  film on a (001) $\text{TiO}_2$  substrate have been recorded using as poles the  $\text{CeO}_2$  (111) reflection ( $2\theta = 28.54^\circ$ ) and the  $\text{TiO}_2$  (111) reflection ( $2\theta = 41.22^\circ$ ) at  $\psi$  tilt angles of  $54.75^\circ$  and  $42.24^\circ$ , respectively. The two



**Figure 7.** Model of cell lattice matching for a (111)  $\text{CeO}_2$  film and (001)  $\text{TiO}_2$  substrate planes. Epitaxial relations are as follows:  $\text{CeO}_2 \langle 111 \rangle \parallel \text{TiO}_2 \langle 001 \rangle$ ,  $\text{CeO}_2 \langle -110 \rangle \parallel \text{TiO}_2 \langle 100 \rangle$ ;  $\text{CeO}_2 \langle 111 \rangle \parallel \text{TiO}_2 \langle 001 \rangle$ ,  $\text{CeO}_2 \langle -110 \rangle \parallel \text{TiO}_2 \langle 010 \rangle$ .

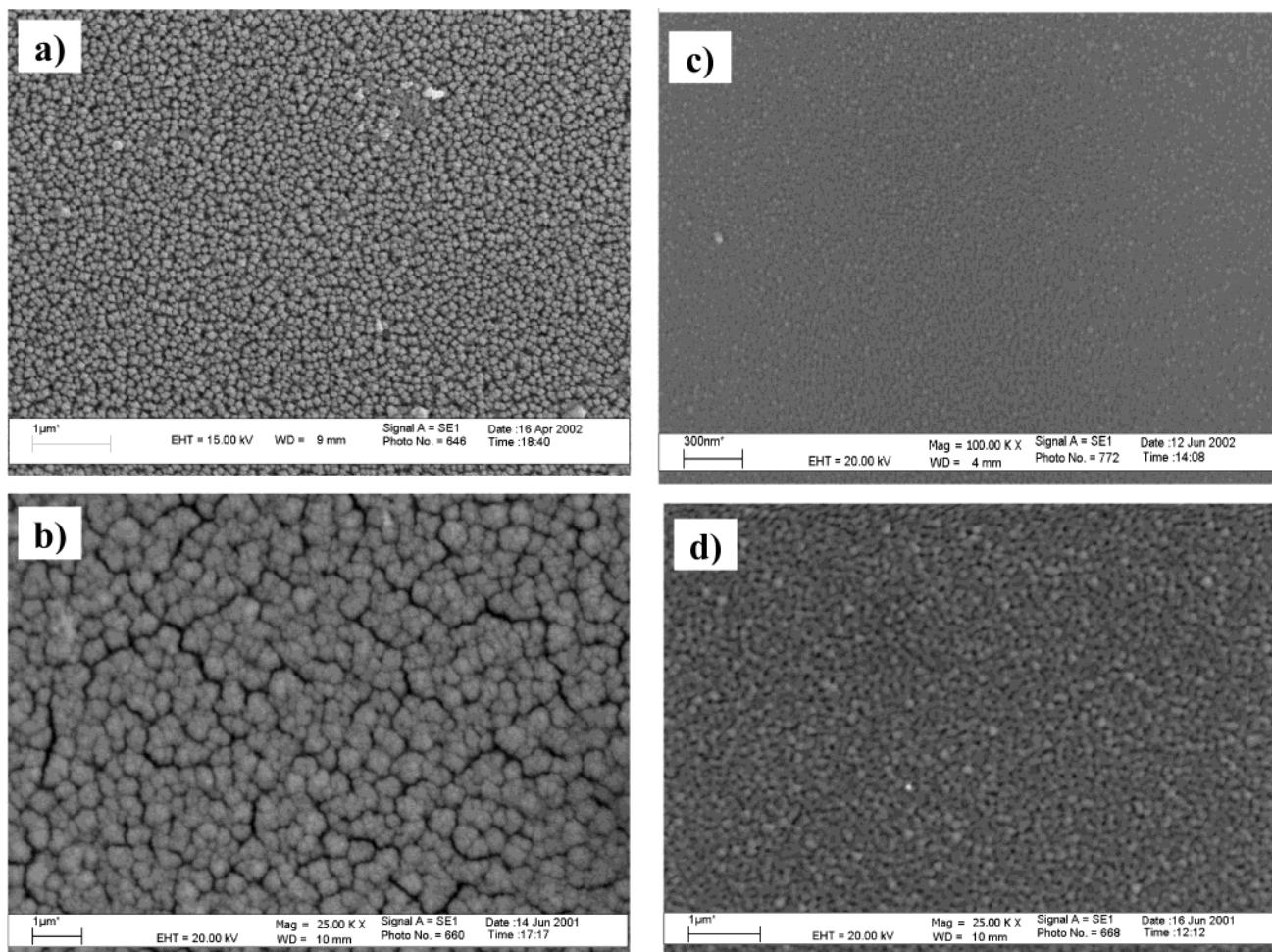


**Figure 8.**  $\Phi$ -scans using the (111) or (200)  $\text{CeO}_2$  reflection as the pole for cerium dioxide films and the (111)  $\text{TiO}_2$  reflection for a rutile substrate: (a)  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  films; (b)  $\langle 001 \rangle$   $\text{TiO}_2$  substrate; (c)  $\langle 111 \rangle$ -oriented  $\text{CeO}_2$  films.

peaks from the  $\text{CeO}_2$  {111} planes are coincident with the {111}  $\text{TiO}_2$   $\Phi$  positions. This clearly demonstrates that the  $\text{CeO}_2$  unit cell is not rotated in the plane with respect to the  $\text{TiO}_2$  substrate; thus, the [100] and [010] axis directions are aligned for the film and the substrate (i.e.,  $\text{CeO}_2 \langle 100 \rangle \parallel \text{TiO}_2 \langle 100 \rangle$ ).

(38) Gorbenko, O.; Graboy, I. E.; Novozhilou, M. A.; Kaul, A. R.; Wahl, G.; Svetchnikov V. L. *J. Phys. IV* **2001**, Pr3/247–Pr3/254.

(39) Molodyk, A. A.; Korsakov, I. E.; Novojilov, M. A.; Graboy, I. E.; Andrey, R. K. *Chem. Vap. Deposition* **2000**, 6, 133–138.



**Figure 9.** SEM images of surface morphologies of  $\langle 100 \rangle$   $\text{CeO}_2$  epitaxial films deposited at (a) 450 °C, (b) 750 °C, (c) 750 °C and slowly cooled, and of  $\langle 111 \rangle$   $\text{CeO}_2$  epitaxial films deposited at 1050 °C (d).

Similarly, the  $\Phi$ -scan of the  $\text{CeO}_2$  (200) peak in a pure  $\langle 111 \rangle$  film is displayed in Figure 8c. Of course, instead of the three peaks expected for the (200) reflection at  $\psi = 54.75^\circ$ , 12 equally spaced ( $30^\circ$ ) peaks are again present.

The in-plane mosaic spread of both  $\langle 100 \rangle$  and  $\langle 111 \rangle$  orientations of  $\text{CeO}_2$  films has been estimated from fwhm values of the  $\Phi$  scan peaks. Differences in the angular spreading are evident for the  $\langle 111 \rangle$  and  $\langle 100 \rangle$  textures. In the former case, the  $\Phi$  scan fwhms are quite small ( $\sim 1.17^\circ$ ) and, in any case, slightly higher than the mosaicity of the films measured by the rocking curves. Broader distributions (close to  $2^\circ$ ) were measured for the  $\langle 100 \rangle$  texturing. The broadening of  $\Delta\Phi$  observed in  $\langle 100 \rangle$ -oriented films, larger than that in  $\langle 111 \rangle$ , can be due to the elastic strain caused by the lattice mismatch. Note, however, that the epitaxial quality of present  $\langle 100 \rangle$   $\text{CeO}_2$  films is comparable to that previously observed for  $\text{CeO}_2$  films deposited by MOCVD on single-crystal substrates<sup>40</sup> having a lattice mismatch greater than 5%.

The cube-on-cube epitaxial growth of the  $\langle 100 \rangle$   $\text{CeO}_2$  films on  $\text{TiO}_2$  occurs despite the large mismatch. Such a difference certainly causes a sizable distortion of the  $\text{CeO}_2$  cell for lattice matching epitaxy. In particular, the  $\langle 100 \rangle$ -oriented films, whose [001] and [010] directions

lie in-plane, have single-unit cell lattice mismatch of 17%. This mismatch, however, can be greatly reduced by matching five unit cells of  $\text{CeO}_2$  to six unit cells of  $\text{TiO}_2$ . In this case the mismatch is 1.8%. Although this model implies a large number of Ce atoms are not coincident with the substrate sites, the small distortion allows this domain epitaxy.

For  $\langle 111 \rangle$ -oriented growth, the in-plane lattice directions are  $[-110]$  and  $[11-2]$ . For the  $[-110]$  direction, simple unit cell matching results in a strain of 50%, which can be reduced to 1.9% by matching one unit cell of  $\text{CeO}_2$  to two unit cells of  $\text{TiO}_2$ . For the  $[11-2]$  direction, a 1:1 lattice ratio results in a mismatch of 17%. However, by matching six unit cells of  $\text{TiO}_2$  to five of  $\text{CeO}_2$  along this direction, the mismatch is only 1.8%.

Simple lattice mismatch considerations could suggest that the  $\langle 111 \rangle$  orientation should be preferred to the  $\langle 100 \rangle$ , due to the slightly smaller lattice mismatch with the substrate and to the lower surface energy for a face-centered cubic (fcc) structure.<sup>41</sup> The reason for the occurrence of both  $\langle 100 \rangle$  and  $\langle 111 \rangle$  orientations of  $\text{CeO}_2$  films is not fully understood. Nevertheless, a suitable rationale can be found in analogy to the case of  $\text{CeO}_2$  films grown on random substrates, such as Hastelloy

(40) Inagaki, T.; Yoshimura, Y.; Kanda, Y.; Matsumoto, Y.; Minami, K. *Nucl. Instrum. Methods A* **2000**, *443*, 126–135.

(41) Schowalter, L. J.; Fathauer, R. W.; Goehner, R. P.; Turner, L. G.; De Blois, R. W.; Hashimoto, S.; Peng, J. L.; Giboon, W. M.; Krusius, J. P. *J. Appl. Phys.* **1985**, *58*, 302–306.

C276 and polycrystalline YSZ.<sup>34</sup> In those cases,  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  films were grown in a homologous temperature ( $T_h = T_{\text{dep}}/T_{\text{melting}}$ ) range<sup>42</sup> in which the fiber texture can be rationalized by a model for texture development based on a subtle balance between differences of surface energies and elastic energies of differently oriented grains.<sup>43</sup> For example, for a fcc film, the  $\langle 111 \rangle$  orientation is favored by surface energy, even though in the mentioned deposition temperature range some modifications to the  $\langle 100 \rangle$  texturing can be induced by the different elastic energy stored in grains with different orientations. At higher deposition temperatures, when crystallites are formed under higher surface mobilities, it is reasonable to assume that lattice match could play a more important role in determining the preferential orientation.

Note, however, that the reason for the occurrence of both  $\langle 100 \rangle$  and  $\langle 111 \rangle$  orientations cannot be fully rationalized, and in addition, there are no other results to be compared with present data since to our knowledge  $\text{CeO}_2$  films have never been grown on  $\text{TiO}_2$ .

**Morphological Characterization.** The film surface has been investigated by scanning electron microscopy (SEM). The  $\text{CeO}_2$  films exhibit a surface structure consisting of grains with different shapes and size depending on both deposition temperatures and different structural orientations. Two grain types can be distinguished: (i) nanostructured grains with rounded shape and (ii) pyramidal caps with a more random shape.

The SEM image of a  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  film, deposited at 450 °C, is shown in Figure 9a. The morphology consists of well-connected pyramidal-shaped grains with a diameter of 100 nm. These regularly shaped grains grow oriented in the  $[100]$  direction and they are capped by pyramids, the four triangular faces being  $\{111\}$  planes. There is some evidence of the same morphology from previous studies on the growth of  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  film on a random Hastelloy C276 substrate.<sup>34</sup>

The SEM image of  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  film deposited at 750 °C shows a cracked surface (Figure 9b). Crack formation may be associated with stress in the  $\text{CeO}_2$  layers. This is possibly due to the lattice mismatch between  $\text{CeO}_2$  films and  $\text{TiO}_2$  substrates and/or to different thermal expansion coefficients. On the other hand, this problem has been overcome by a slow cooling (3 °C/min) of  $\text{CeO}_2$  films under vacuum. The surface morphology of a slowly cooled sample is shown in Figure 9c. The nanostructured surface consists of few tens of nanometer wide grains, forming a smooth and homogeneous surface suitable for the growth of superconducting thin films.

(42) Machlin, E. S. *Materials Science in Microelectronics*; Gyro: New York, 1995.

(43) Bunshah, R. F.; Blocher, J. M., Jr.; Mattox, D. M. *Deposition Technologies for Films and Coatings*; Noesy Publication: Park Ridge, NY, 1982.

In Figure 9d is reported the SEM image of a  $\langle 111 \rangle$ -oriented film deposited at 1050 °C. The sample presents a different morphology which consists of a fine-grained structure and rounded grains having diameters of about 100 nm. The grains form a regularly and tightly packed array, indicating a good in-plane and out-of-plane alignment.

Finally, it is likely that there is no influence of the film thickness, since it varies only from 0.7 to 1.1  $\mu\text{m}$ , on the surface morphology.

The chemical purity of the samples has been determined by wavelength-dispersive X-ray (WDX) analyses. All the experiments on present films indicate negligible C (0.2%) and F (0.1%) contamination. The WDX detection limits are about 200 and 400 ppm, respectively, under the conditions reported in the Experimental Section.

## Conclusions

In the present study we report results of cerium dioxide epitaxial deposition on rutile as a potential buffer layer for superconducting films growth. The significant lattice mismatch (about 17%) between a  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  film and  $\text{TiO}_2$  (001) substrate points to a critical epitaxial growth. Present data, however, show that epitaxial  $\langle 100 \rangle$   $\text{CeO}_2$  films can be grown on (001)-oriented  $\text{TiO}_2$  substrates by the MOCVD technique.

It has been found that texturing and epitaxy of  $\text{CeO}_2$  films is very sensitive to the deposition temperature. The epitaxial growth of  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  films has been tailored by varying deposition parameters and monitored by a complete structural characterization.  $\theta$ – $2\theta$  scans and X-ray pole figures show that  $\langle 100 \rangle$   $\text{CeO}_2$  films with good crystallinity can be epitaxially deposited on (001)  $\text{TiO}_2$  substrates in the 450–750 °C temperature range.

Detailed information has been obtained from the related rocking curves whose fwhms indicate that higher temperature values improve the out-of-plane alignment. The  $\langle 100 \rangle$ -oriented  $\text{CeO}_2$  films grow cube-on-cube on the  $\text{TiO}_2$  substrate, that is,  $\text{CeO}_2 \langle 100 \rangle \parallel \text{TiO}_2 \langle 100 \rangle$ .

Moreover,  $\text{CeO}_2$  films, grown at deposition temperatures higher than 750 °C, show a  $\langle 111 \rangle$  orientation. Pole figures of the (200) reflection for  $\langle 111 \rangle$   $\text{CeO}_2$  films show 12 maxima, indicating the presence of at least two in-plane variants.

The evolution of fiber texture, upon varying the growth conditions, has been interpreted in terms of surface mobility and lattice matching.

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